

SM_Kerr

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1 Kerr spacetime

This notebook demonstrates a few capabilities of SageMath in computations regarding Kerr spacetime. The corresponding tools have been developed within the [SageManifolds](#) project.

Click [here](#) to download the notebook file (ipynb format). To run it, you must start SageMath within the Jupyter notebook, via the command `sage -n jupyter`

NB: a version of SageMath at least equal to 8.2 is required to run this notebook:

```
[1]: version()
```

```
[1]: 'SageMath version 9.2.beta13, Release Date: 2020-09-21'
```

First we set up the notebook to display mathematical objects using LaTeX rendering:

```
[2]: %display latex
```

and we initialize a time counter for benchmarking:

```
[3]: import time
     comput_time0 = time.perf_counter()
```

Since some computations are quite heavy, we ask for running them in parallel on 8 threads:

```
[4]: Parallelism().set(nproc=8)
```

1.1 Spacetime manifold

We declare the Kerr spacetime (or more precisely the part of it covered by Boyer-Lindquist coordinates) as a 4-dimensional Lorentzian manifold \mathcal{M} :

```
[5]: M = Manifold(4, 'M', latex_name=r'\mathcal{M}', structure='Lorentzian')
     print(M)
```

4-dimensional Lorentzian manifold M

We then introduce the standard Boyer-Lindquist coordinates as a chart BL (for *Boyer-Lindquist*) on \mathcal{M} , via the method `chart()`, the argument of which is a string (delimited by `r"..."` because of

the backslash symbols) expressing the coordinates names, their ranges (the default is $(-\infty, +\infty)$) and their LaTeX symbols:

```
[6]: BL.<t,r,th,ph> = M.chart(r"t r th:(0,pi):\theta ph:(0,2*pi):\phi")
print(BL); BL
```

Chart (M, (t, r, th, ph))

```
[6]: (M, (t, r, theta, phi))
```

```
[7]: BL[0], BL[1]
```

```
[7]: (t, r)
```

Metric tensor

The 2 parameters m and a of the Kerr spacetime are declared as symbolic variables:

```
[8]: var('m, a', domain='real')
```

```
[8]: (m, a)
```

We get the (yet undefined) spacetime metric:

```
[9]: g = M.metric()
```

The metric is set by its components in the coordinate frame associated with Boyer-Lindquist coordinates, which is the current manifold's default frame:

```
[10]: rho2 = r^2 + (a*cos(th))^2
Delta = r^2 - 2*m*r + a^2
g[0,0] = -(1-2*m*r/rho2)
g[0,3] = -2*a*m*r*sin(th)^2/rho2
g[1,1], g[2,2] = rho2/Delta, rho2
g[3,3] = (r^2+a^2+2*m*r*(a*sin(th))^2/rho2)*sin(th)^2
g.display()
```

```
[10]: g = ( (frac(2*m*r, a^2*cos(theta)^2+r^2) - 1) dt \otimes dt + ( -frac(2*a*m*r*sin(theta)^2, a^2*cos(theta)^2+r^2) dt \otimes dphi + (frac(a^2*cos(theta)^2+r^2, a^2-2*m*r+r^2) dr \otimes dr + (a^2*cos(theta)^2+r^2) dtheta \otimes dtheta + ( -frac(2*a*m*r*sin(theta)^2, a^2*cos(theta)^2+r^2) dphi \otimes dt + (frac(2*a^2*m*r*sin(theta)^2, a^2*cos(theta)^2+r^2) + a^2 + r^2) sin(theta)^2 dphi \otimes dphi
```

A matrix view of the components with respect to the manifold's default vector frame:

```
[11]: g[:]
```

```
[11]: ( (frac(2*m*r, a^2*cos(theta)^2+r^2) - 1, 0, 0, -frac(2*a*m*r*sin(theta)^2, a^2*cos(theta)^2+r^2)
0, frac(a^2*cos(theta)^2+r^2, a^2-2*m*r+r^2), 0, 0
0, 0, a^2*cos(theta)^2+r^2, 0
-fraction(2*a*m*r*sin(theta)^2, a^2*cos(theta)^2+r^2), 0, 0, (frac(2*a^2*m*r*sin(theta)^2, a^2*cos(theta)^2+r^2) + a^2 + r^2) sin(theta)^2 )
```

The list of the non-vanishing components:

```
[12]: g.display_comp()
```

```
[12]:
```

$$\begin{aligned}g_{tt} &= \frac{2mr}{a^2 \cos(\theta)^2 + r^2} - 1 \\g_{t\phi} &= -\frac{2amr \sin(\theta)^2}{a^2 \cos(\theta)^2 + r^2} \\g_{rr} &= \frac{a^2 \cos(\theta)^2 + r^2}{a^2 - 2mr + r^2} \\g_{\theta\theta} &= a^2 \cos(\theta)^2 + r^2 \\g_{\phi t} &= -\frac{2amr \sin(\theta)^2}{a^2 \cos(\theta)^2 + r^2} \\g_{\phi\phi} &= \left(\frac{2a^2mr \sin(\theta)^2}{a^2 \cos(\theta)^2 + r^2} + a^2 + r^2 \right) \sin(\theta)^2\end{aligned}$$

Levi-Civita Connection

The Levi-Civita connection ∇ associated with g :

```
[13]: nabra = g.connection()
print(nabra)
```

Levi-Civita connection `nabra_g` associated with the Lorentzian metric g on the 4-dimensional Lorentzian manifold M

Let us verify that the covariant derivative of g with respect to ∇ vanishes identically:

```
[14]: nabra(g) == 0
```

```
[14]: True
```

Another view of the above property:

```
[15]: nabra(g).display()
```

```
[15]:  $\nabla_g g = 0$ 
```

The nonzero Christoffel symbols (skipping those that can be deduced by symmetry of the last two indices):

```
[16]: g.christoffel_symbols_display()
```

```
[16]:
```

$$\begin{aligned}
\Gamma^t_{tr} &= -\frac{a^4 m - m r^4 - (a^4 m + a^2 m r^2) \sin(\theta)^2}{a^2 r^4 - 2 m r^5 + r^6 + (a^6 - 2 a^4 m r + a^4 r^2) \cos(\theta)^4 + 2 (a^4 r^2 - 2 a^2 m r^3 + a^2 r^4) \cos(\theta)^2} \\
\Gamma^t_{t\theta} &= -\frac{2 a^2 m r \cos(\theta) \sin(\theta)}{a^4 \cos(\theta)^4 + 2 a^2 r^2 \cos(\theta)^2 + r^4} \\
\Gamma^t_{r\phi} &= -\frac{(a^3 m r^2 + 3 a m r^4 - (a^5 m - a^3 m r^2) \cos(\theta)^2) \sin(\theta)^2}{a^2 r^4 - 2 m r^5 + r^6 + (a^6 - 2 a^4 m r + a^4 r^2) \cos(\theta)^4 + 2 (a^4 r^2 - 2 a^2 m r^3 + a^2 r^4) \cos(\theta)^2} \\
\Gamma^t_{\theta\phi} &= -\frac{2 (a^5 m r \cos(\theta) \sin(\theta)^5 - (a^5 m r + a^3 m r^3) \cos(\theta) \sin(\theta)^3)}{a^6 \cos(\theta)^6 + 3 a^4 r^2 \cos(\theta)^4 + 3 a^2 r^4 \cos(\theta)^2 + r^6} \\
\Gamma^r_{tt} &= \frac{a^2 m r^2 - 2 m^2 r^3 + m r^4 - (a^4 m - 2 a^2 m^2 r + a^2 m r^2) \cos(\theta)^2}{a^6 \cos(\theta)^6 + 3 a^4 r^2 \cos(\theta)^4 + 3 a^2 r^4 \cos(\theta)^2 + r^6} \\
\Gamma^r_{t\phi} &= -\frac{(a^3 m r^2 - 2 a m^2 r^3 + a m r^4 - (a^5 m - 2 a^3 m^2 r + a^3 m r^2) \cos(\theta)^2) \sin(\theta)^2}{a^6 \cos(\theta)^6 + 3 a^4 r^2 \cos(\theta)^4 + 3 a^2 r^4 \cos(\theta)^2 + r^6} \\
\Gamma^r_{rr} &= \frac{a^2 r - m r^2 + (a^2 m - a^2 r) \cos(\theta)^2}{a^2 r^2 - 2 m r^3 + r^4 + (a^4 - 2 a^2 m r + a^2 r^2) \cos(\theta)^2} \\
\Gamma^r_{r\theta} &= -\frac{a^2 \cos(\theta) \sin(\theta)}{a^2 \cos(\theta)^2 + r^2} \\
\Gamma^r_{\theta\theta} &= -\frac{a^2 r - 2 m r^2 + r^3}{a^2 \cos(\theta)^2 + r^2} \\
\Gamma^r_{\phi\phi} &= \frac{(a^4 m r^2 - 2 a^2 m^2 r^3 + a^2 m r^4 - (a^6 m - 2 a^4 m^2 r + a^4 m r^2) \cos(\theta)^2) \sin(\theta)^4 - (a^2 r^5 - 2 m r^6 + r^7 + (a^6 r - 2 a^4 m r^2 + a^4 r^3) \cos(\theta)^4 + 2 (a^4 r^3 - a^6 \cos(\theta)^6 + 3 a^4 r^2 \cos(\theta)^4 + 3 a^2 r^4 \cos(\theta)^2 + r^6)) \sin(\theta)^2}{a^6 \cos(\theta)^6 + 3 a^4 r^2 \cos(\theta)^4 + 3 a^2 r^4 \cos(\theta)^2 + r^6} \\
\Gamma^\theta_{tt} &= -\frac{2 a^2 m r \cos(\theta) \sin(\theta)}{a^6 \cos(\theta)^6 + 3 a^4 r^2 \cos(\theta)^4 + 3 a^2 r^4 \cos(\theta)^2 + r^6} \\
\Gamma^\theta_{t\phi} &= \frac{2 (a^3 m r + a m r^3) \cos(\theta) \sin(\theta)}{a^6 \cos(\theta)^6 + 3 a^4 r^2 \cos(\theta)^4 + 3 a^2 r^4 \cos(\theta)^2 + r^6} \\
\Gamma^\theta_{rr} &= \frac{a^2 \cos(\theta) \sin(\theta)}{a^2 r^2 - 2 m r^3 + r^4 + (a^4 - 2 a^2 m r + a^2 r^2) \cos(\theta)^2} \\
\Gamma^\theta_{r\theta} &= \frac{r}{a^2 \cos(\theta)^2 + r^2} \\
\Gamma^\theta_{\theta\theta} &= -\frac{a^2 \cos(\theta) \sin(\theta)}{a^2 \cos(\theta)^2 + r^2} \\
\Gamma^\theta_{\phi\phi} &= -\frac{((a^6 - 2 a^4 m r + a^4 r^2) \cos(\theta)^5 + 2 (a^4 r^2 - 2 a^2 m r^3 + a^2 r^4) \cos(\theta)^3 + (2 a^4 m r + 4 a^2 m r^3 + a^2 r^4 + r^6) \cos(\theta)) \sin(\theta)}{a^6 \cos(\theta)^6 + 3 a^4 r^2 \cos(\theta)^4 + 3 a^2 r^4 \cos(\theta)^2 + r^6} \\
\Gamma^\phi_{tr} &= -\frac{a^3 m \cos(\theta)^2 - a m r^2}{a^2 r^4 - 2 m r^5 + r^6 + (a^6 - 2 a^4 m r + a^4 r^2) \cos(\theta)^4 + 2 (a^4 r^2 - 2 a^2 m r^3 + a^2 r^4) \cos(\theta)^2} \\
\Gamma^\phi_{t\theta} &= -\frac{2 a m r \cos(\theta)}{(a^4 \cos(\theta)^4 + 2 a^2 r^2 \cos(\theta)^2 + r^4) \sin(\theta)} \\
\Gamma^\phi_{r\phi} &= -\frac{a^2 m r^2 + 2 m r^4 - r^5 + (a^4 m - a^4 r) \cos(\theta)^4 - (a^4 m - a^2 m r^2 + 2 a^2 r^3) \cos(\theta)^2}{a^2 r^4 - 2 m r^5 + r^6 + (a^6 - 2 a^4 m r + a^4 r^2) \cos(\theta)^4 + 2 (a^4 r^2 - 2 a^2 m r^3 + a^2 r^4) \cos(\theta)^2} \\
\Gamma^\phi_{\theta\phi} &= \frac{a^4 \cos(\theta) \sin(\theta)^4 - 2 (a^4 - a^2 m r + a^2 r^2) \cos(\theta) \sin(\theta)^2 + (a^4 + 2 a^2 r^2 + r^4) \cos(\theta)}{(a^4 \cos(\theta)^4 + 2 a^2 r^2 \cos(\theta)^2 + r^4) \sin(\theta)}
\end{aligned}$$

Killing vectors

The default vector frame on the spacetime manifold is the coordinate basis associated with Boyer-Lindquist coordinates:

```
[17]: M.default_frame() is BL.frame()
```

```
[17]: True
```

```
[18]: BL.frame()
```

```
[18]: (M, (∂/∂t, ∂/∂r, ∂/∂θ, ∂/∂φ))
```

Let us consider the first vector field of this frame:

```
[19]: xi = BL.frame()[0]; xi
```

```
[19]: ∂/∂t
```

```
[20]: print(xi)
```

Vector field d/dt on the 4-dimensional Lorentzian manifold M

The 1-form associated to it by metric duality is

```
[21]: xi_form = xi.down(g)
xi_form.display()
```

```
[21]:  $\left(\frac{2mr}{a^2 \cos(\theta)^2 + r^2} - 1\right) dt + \left(-\frac{2amr \sin(\theta)^2}{a^2 \cos(\theta)^2 + r^2}\right) d\phi$ 
```

Its covariant derivative is

```
[22]: nab_xi = nabla(xi_form)
print(nab_xi)
nab_xi.display()
```

Tensor field of type $(0,2)$ on the 4-dimensional Lorentzian manifold M

```
[22]:  $\left(\frac{a^2 m \cos(\theta)^2 - mr^2}{a^4 \cos(\theta)^4 + 2a^2 r^2 \cos(\theta)^2 + r^4}\right) dt \otimes dr + \left(\frac{2a^2 mr \cos(\theta) \sin(\theta)}{a^4 \cos(\theta)^4 + 2a^2 r^2 \cos(\theta)^2 + r^4}\right) dt \otimes d\theta +$   
 $\left(-\frac{a^2 m \cos(\theta)^2 - mr^2}{a^4 \cos(\theta)^4 + 2a^2 r^2 \cos(\theta)^2 + r^4}\right) dr \otimes dt + \left(\frac{(a^3 m \cos(\theta)^2 - amr^2) \sin(\theta)^2}{a^4 \cos(\theta)^4 + 2a^2 r^2 \cos(\theta)^2 + r^4}\right) dr \otimes d\phi +$   
 $\left(-\frac{2a^2 mr \cos(\theta) \sin(\theta)}{a^4 \cos(\theta)^4 + 2a^2 r^2 \cos(\theta)^2 + r^4}\right) d\theta \otimes dt + \left(\frac{2(a^3 mr + amr^3) \cos(\theta) \sin(\theta)}{a^4 \cos(\theta)^4 + 2a^2 r^2 \cos(\theta)^2 + r^4}\right) d\theta \otimes d\phi +$   
 $\left(-\frac{(a^3 m \cos(\theta)^2 - amr^2) \sin(\theta)^2}{a^4 \cos(\theta)^4 + 2a^2 r^2 \cos(\theta)^2 + r^4}\right) d\phi \otimes dr + \left(-\frac{2(a^3 mr + amr^3) \cos(\theta) \sin(\theta)}{a^4 \cos(\theta)^4 + 2a^2 r^2 \cos(\theta)^2 + r^4}\right) d\phi \otimes d\theta$ 
```

Let us check that the Killing equation is satisfied:

```
[23]: nab_xi.symmetrize() == 0
```

```
[23]: True
```

Similarly, let us check that $\frac{\partial}{\partial \phi}$ is a Killing vector:

```
[24]: chi = BL.frame()[3] ; chi
```

```
[24]:  $\frac{\partial}{\partial \phi}$ 
```

```
[25]: nabla(chi.down(g)).symmetrize() == 0
```

```
[25]: True
```

Curvature

The Ricci tensor associated with g :

```
[26]: Ric = g.ricci()
print(Ric)
```

Field of symmetric bilinear forms $\text{Ric}(g)$ on the 4-dimensional Lorentzian manifold M

Tensor field of type (4,0) on the 4-dimensional Lorentzian manifold M

The Kretschmann scalar $K := R^{abcd}R_{abcd}$:

```
[37]: Kr_scalar = uR['^{abcd}']*dR['_{abcd}']
      Kr_scalar.display()
```

[37]: $\mathcal{M} \rightarrow \mathbb{R}$
 $(t, r, \theta, \phi) \mapsto -\frac{48(a^6 m^2 \cos(\theta)^6 - 15 a^4 m^2 r^2 \cos(\theta)^4 + 15 a^2 m^2 r^4 \cos(\theta)^2 - m^2 r^6)}{a^{12} \cos(\theta)^{12} + 6 a^{10} r^2 \cos(\theta)^{10} + 15 a^8 r^4 \cos(\theta)^8 + 20 a^6 r^6 \cos(\theta)^6 + 15 a^4 r^8 \cos(\theta)^4 + 6 a^2 r^{10} \cos(\theta)^2 + r^{12}}$

A variant of this expression can be obtained by invoking the `factor()` method on the coordinate function representing the scalar field in the manifold's default chart:

```
[38]: Kr = Kr_scalar.coord_function()
      Kr.factor()
```

[38]:
$$\frac{48(a^2 \cos(\theta)^2 + 4 a r \cos(\theta) + r^2)(a^2 \cos(\theta)^2 - 4 a r \cos(\theta) + r^2)(a \cos(\theta) + r)(a \cos(\theta) - r)m^2}{(a^2 \cos(\theta)^2 + r^2)^6}$$

As a check, we can compare Kr to the formula given by R. Conn Henry, *Astrophys. J.* 535, 350 (2000):

```
[39]: Kr == 48*m^2*(r^6 - 15*r^4*(a*cos(th))^2 + 15*r^2*(a*cos(th))^4
      - (a*cos(th))^6) / (r^2+(a*cos(th))^2)^6
```

[39]: True

The Schwarzschild value of the Kretschmann scalar is recovered by setting $a = 0$:

```
[40]: Kr.expr().subs(a=0)
```

[40]: $\frac{48 m^2}{r^6}$

Let us plot the Kretschmann scalar for $m = 1$ and $a = 0.9$:

```
[41]: K1 = Kr.expr().subs(m=1, a=0.9)
      plot3d(K1, (r, 1, 3), (th, 0, pi), axes_labels=['r', 'theta', 'Kr'])
```

[41]: Graphics3d Object

```
[42]: print("Total elapsed time: {} s".format(time.perf_counter() - comput_time0))
```

Total elapsed time: 483.1138785999992 s